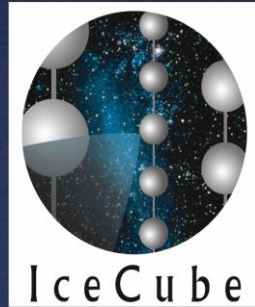


Sterile Neutrinos and IceCube



Warren Huelsnitz
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whuelsn@fnal.gov

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12 May 2011

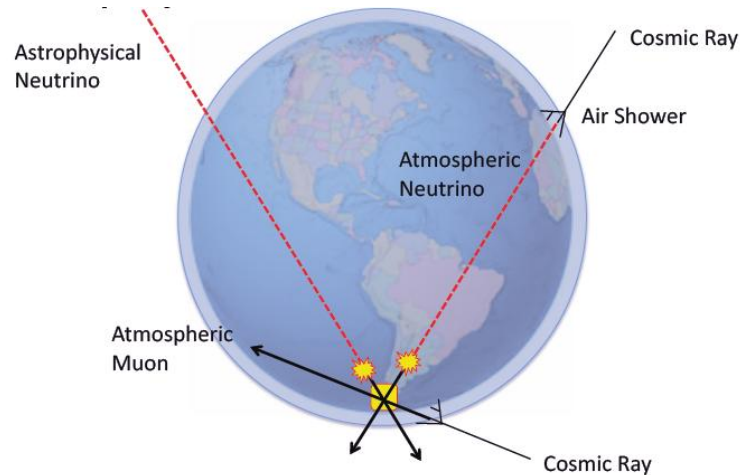
Outline

- Introduction to IceCube
- Atmospheric neutrinos in IceCube
- Sterile neutrinos in IceCube



ScienceDaily (Dec. 19, 2010) — Culminating a decade of planning, innovation and testing, construction of the world's largest neutrino observatory, installed in the ice of the Antarctic plateau at the geographic South Pole, was successfully completed December 18, 2010, New Zealand time.

New window on the universe

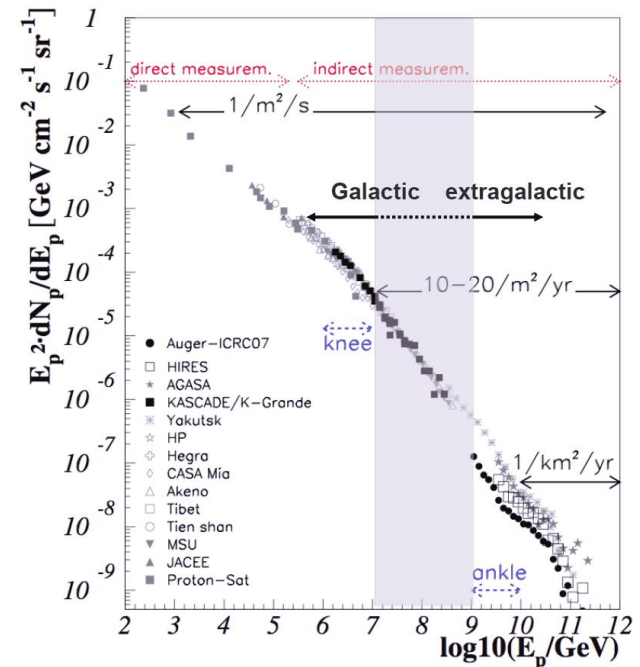
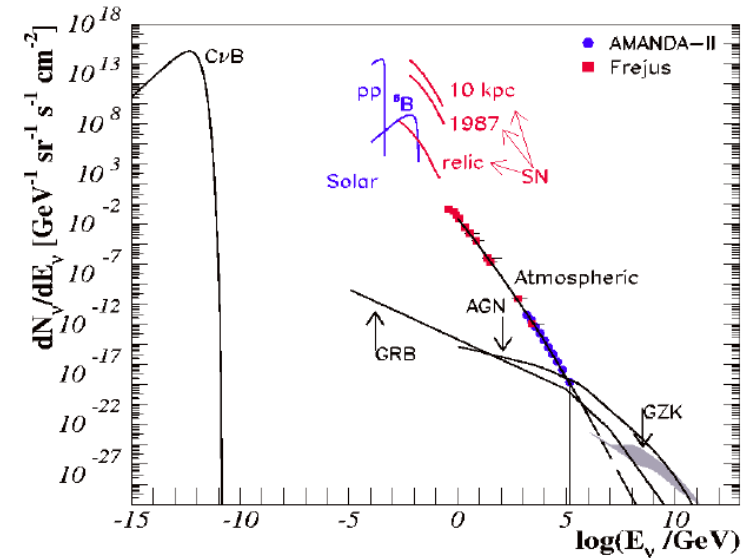


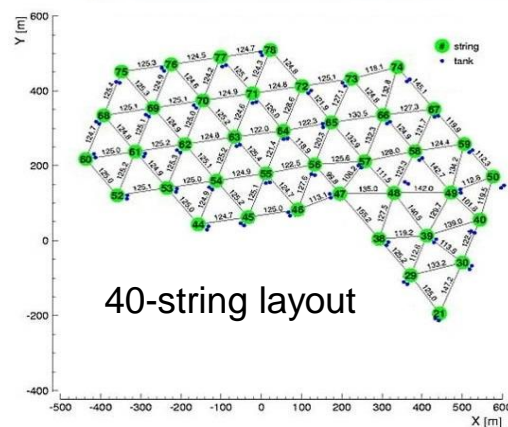
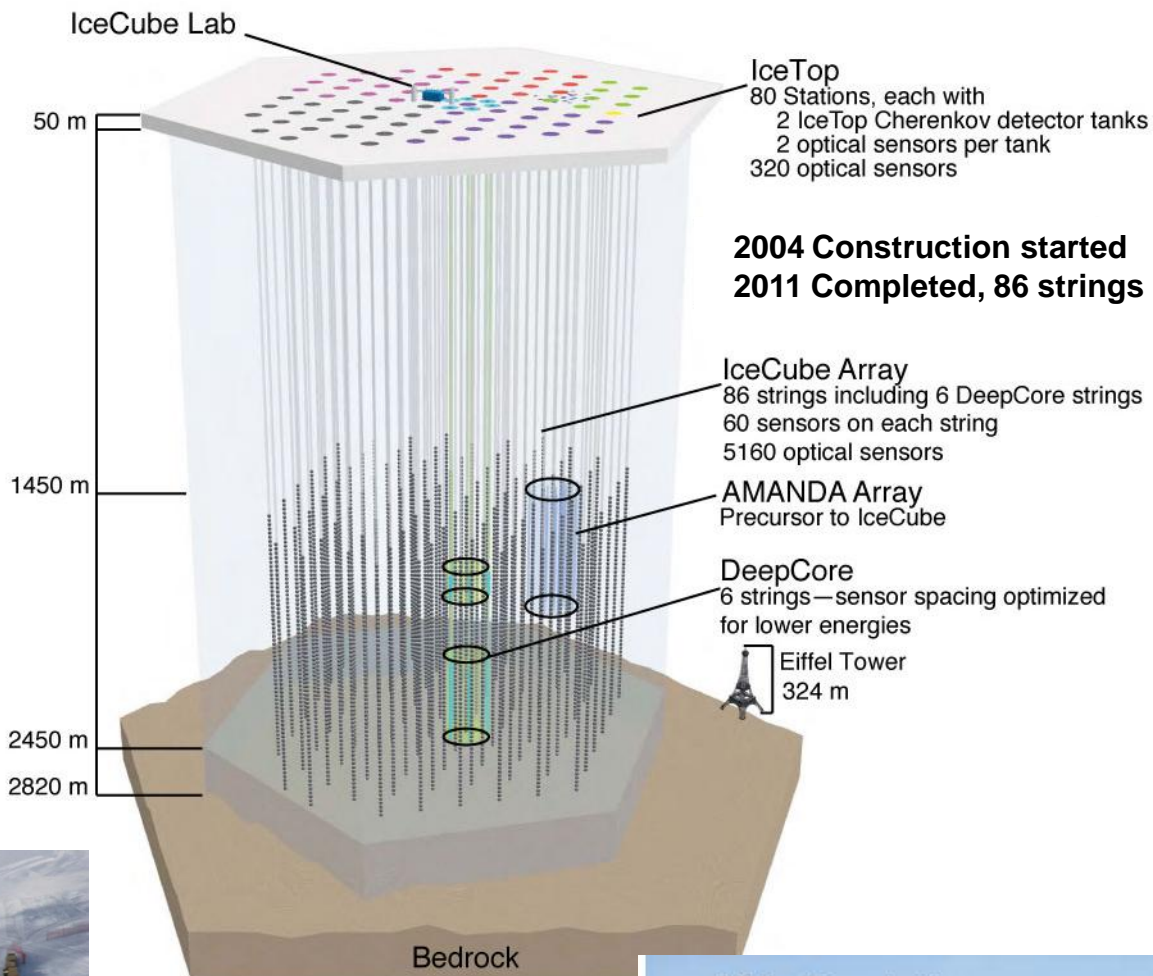
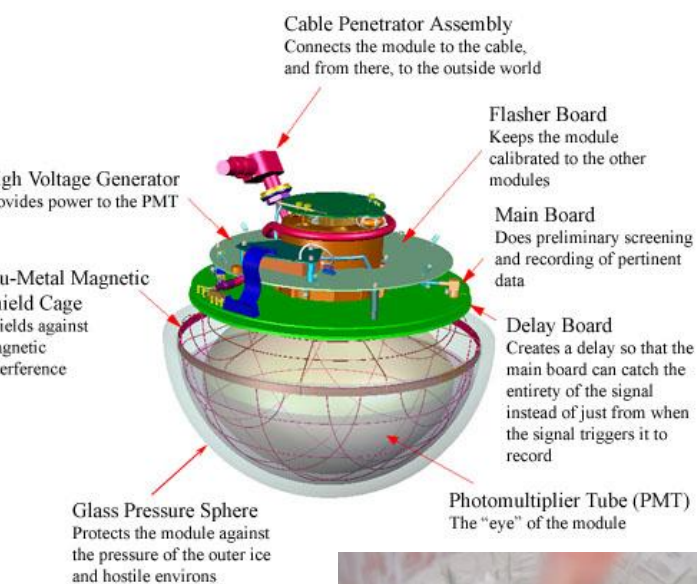
• Why neutrinos?

- No strong or EM interactions
- High Lorentz factors
- Can provide new information about objects and processes in the universe
- Can probe new physics phenomena

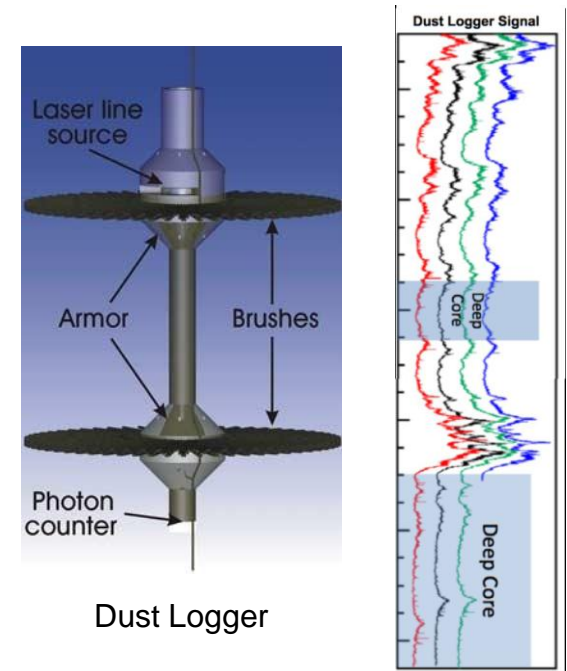
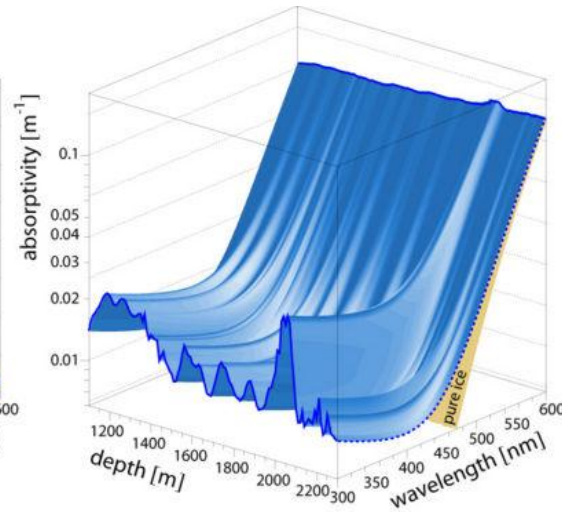
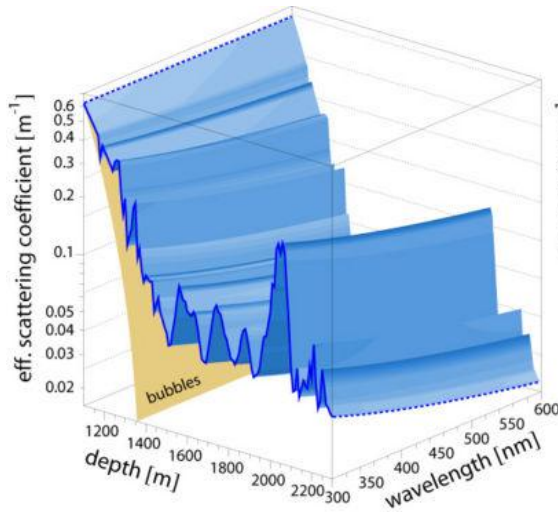
• Why the South Pole?

- Ice provides stable, transparent medium for propagation of Cherenkov photons
- Provides overburden to reduce background of atmospheric muons
- Large target volume for sufficient event rate for astrophysical neutrinos





Ice Properties

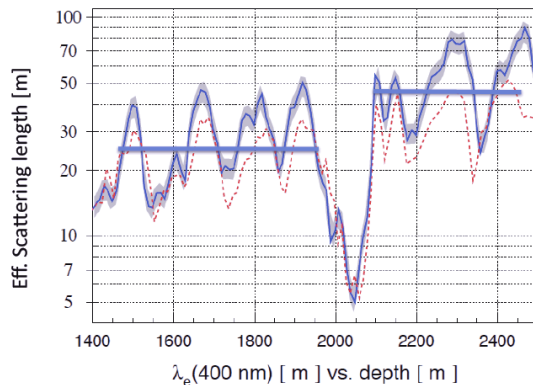


Average optical parameters at 400 nm:

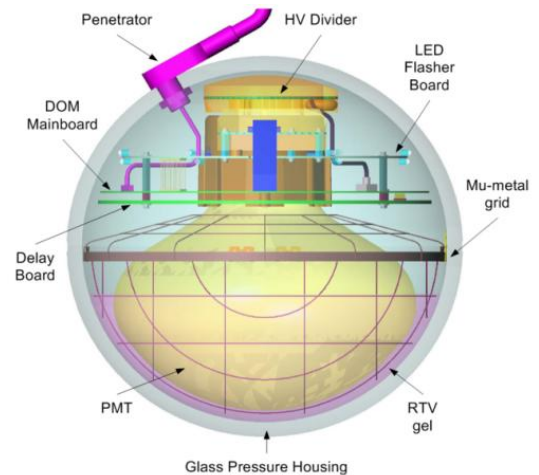
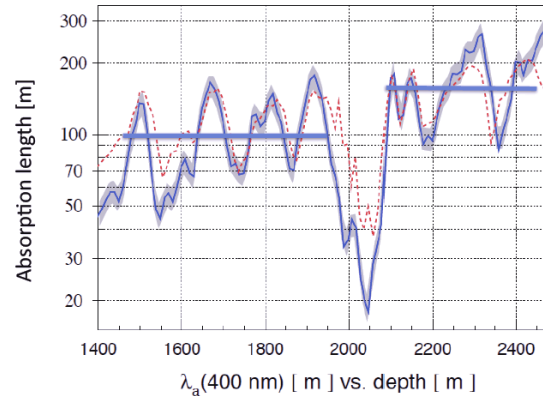
$\lambda_{\text{abs}} \sim 110 \text{ m}$, $\lambda_{\text{sca}} \sim 20 \text{ m}$ above the dust layer

$\lambda_{\text{abs}} \sim 220 \text{ m}$, $\lambda_{\text{sca}} \sim 40 \text{ m}$ below the dust layer

Effective scattering length vs Depth

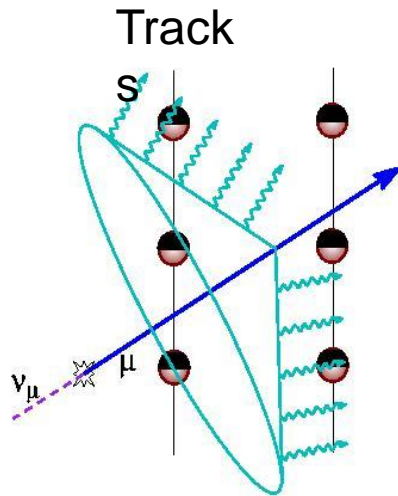


Absorption length vs Depth



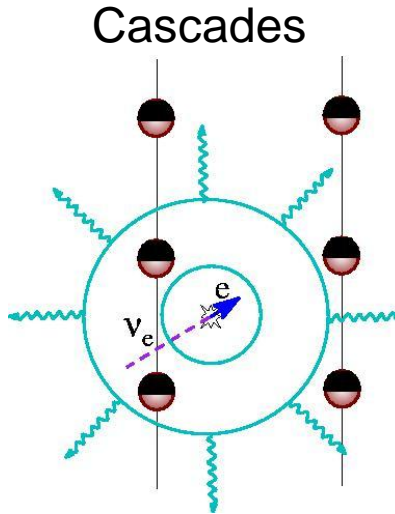
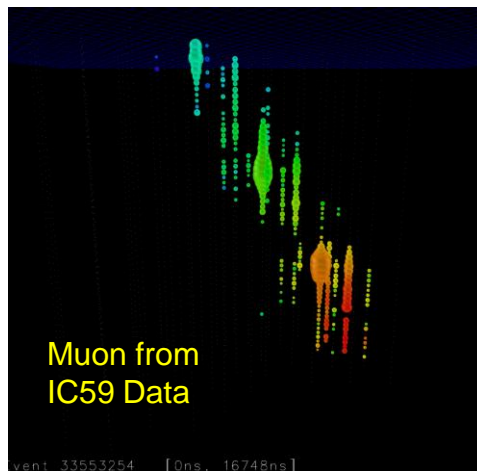
Digital Optical Module (DOM)

Neutrino Detection

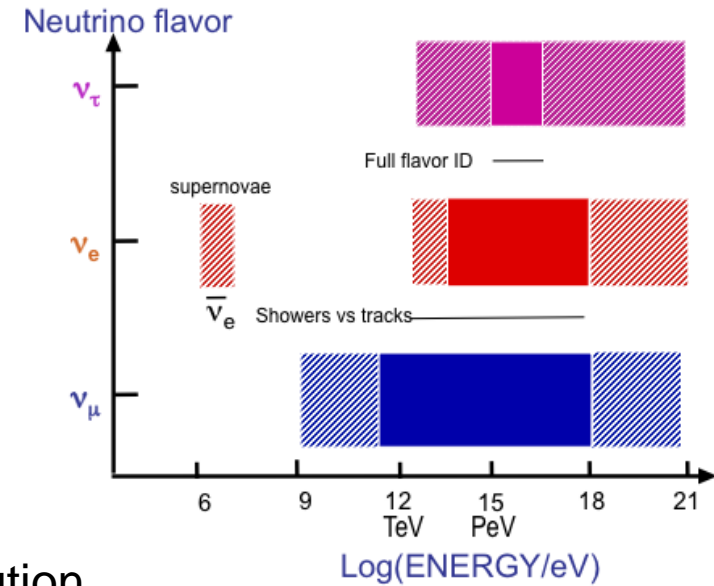
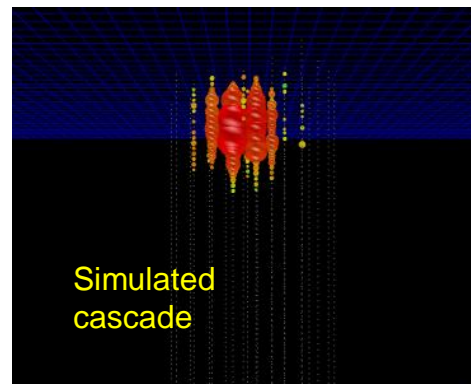


O(km) muon tracks

- ✓ Good angular resolution
- ✓ ν_μ CC events

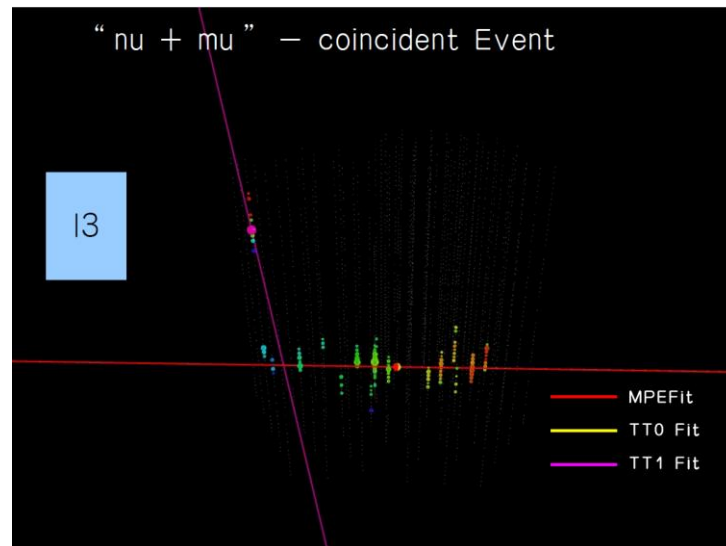
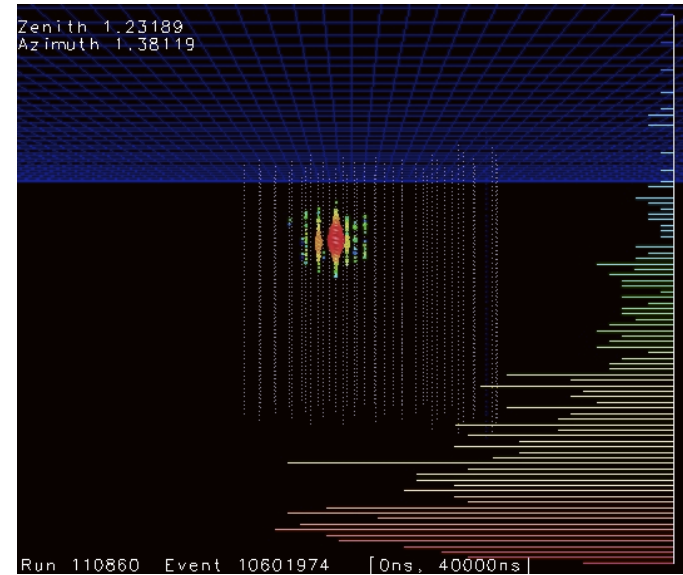
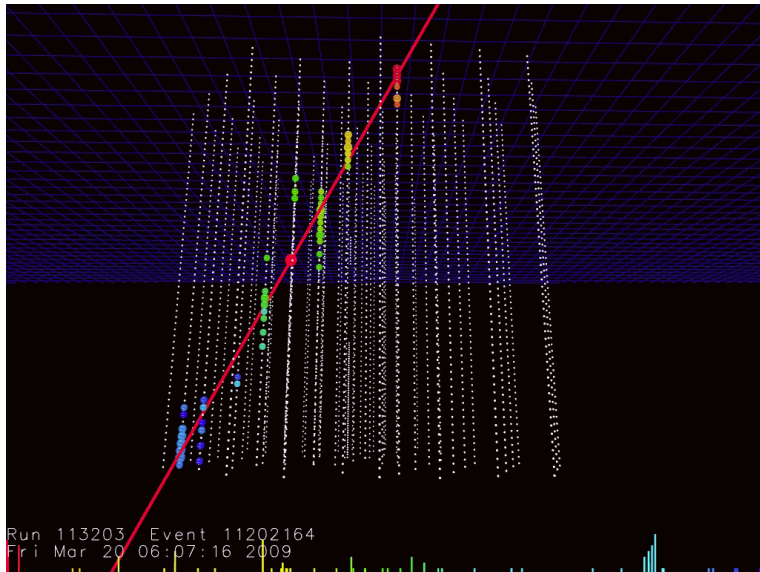


- ✓ Good energy resolution
- ✓ Little directionality
- ✓ 4π sensitivity
- ✓ All flavor sensitivity



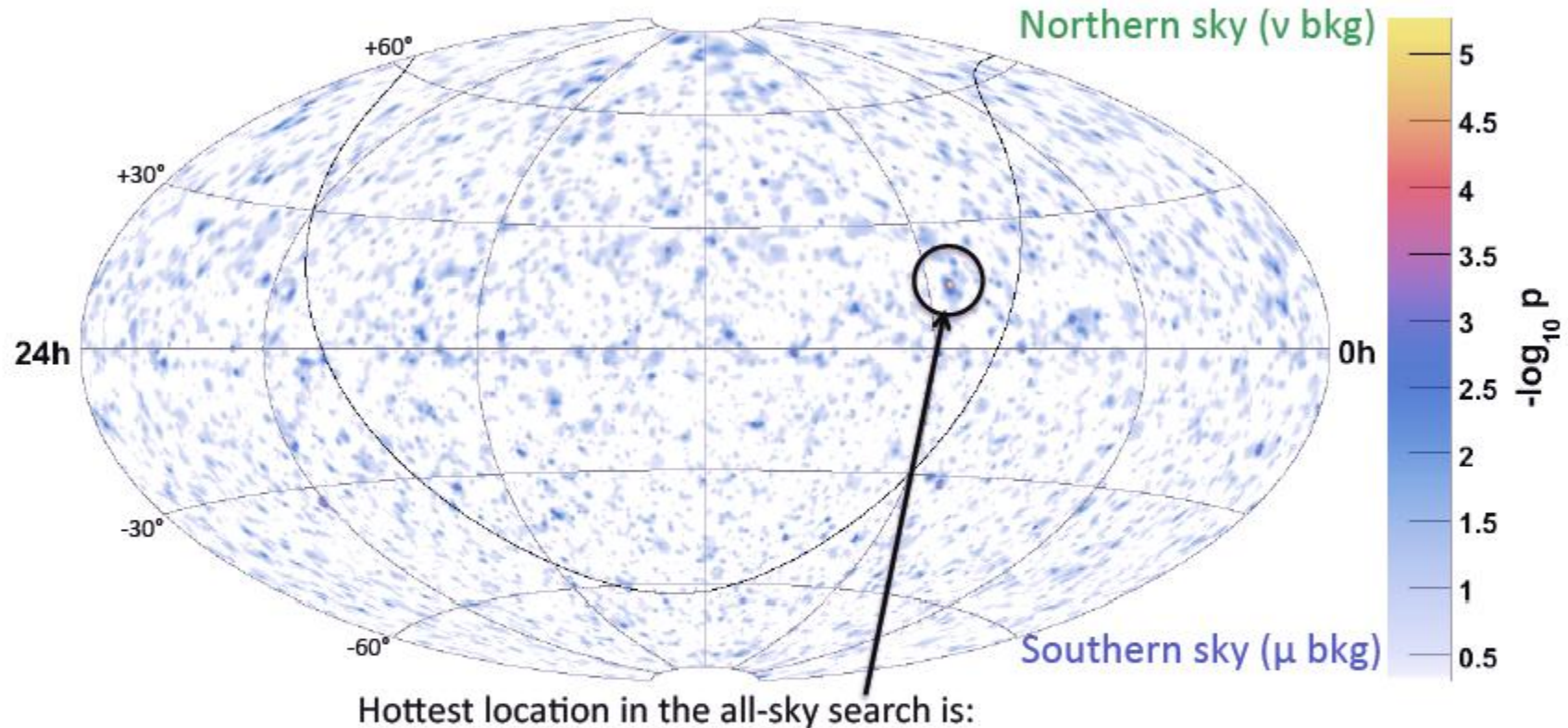
special case:

high energy ν_τ



Astrophysical Point Source Search

IC40 All-Sky Point source search: Likelihood approach



Livetime = 375.5 days

Events = 36900
(14121 up-going,
22779 down-going)

Ra=113.75°, Dec=15.15°

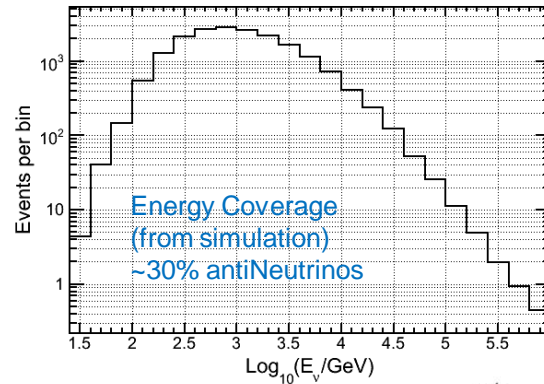
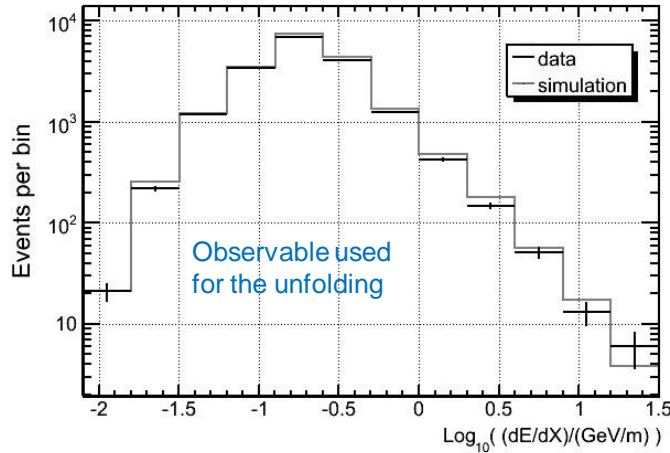
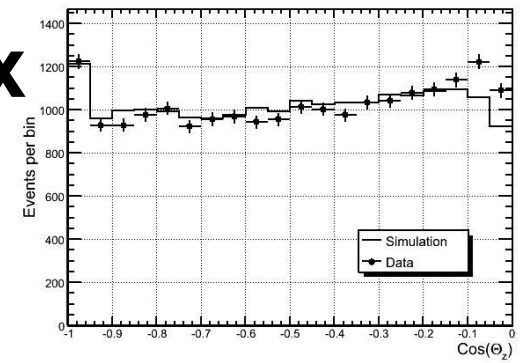
In background simulated trials (scrambling data in r.a.), an equal or greater excess occurs by chance 1817/10000 times.
=> all-sky final p-value is 18%, not significant

arxiv:1012.2137, accepted by Astrophysical Journal

Atmospheric Muon Neutrino Flux

- SVD unfolding method described in A. Höcker and V. Kartvelishvili, [NIM A 372 \(1996\) 469](#)
- Implemented in RooUnfold: <http://hepunix.rl.ac.uk/~adye/software/unfold/RooUnfold.html>

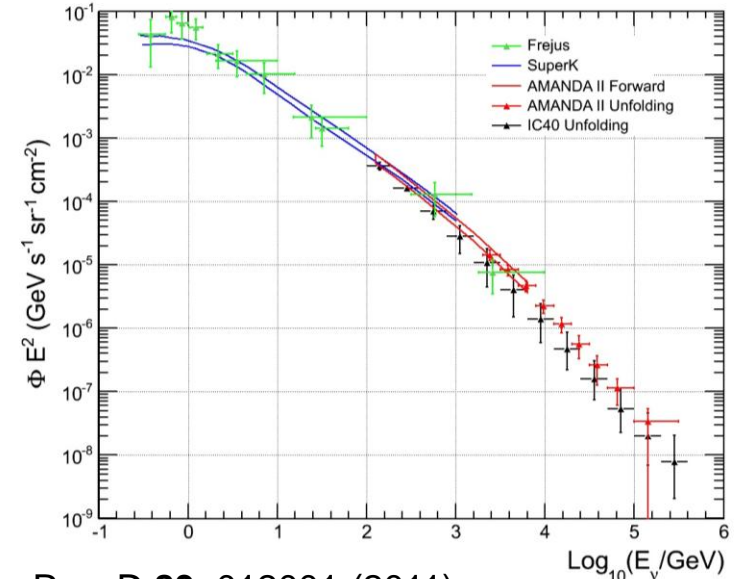
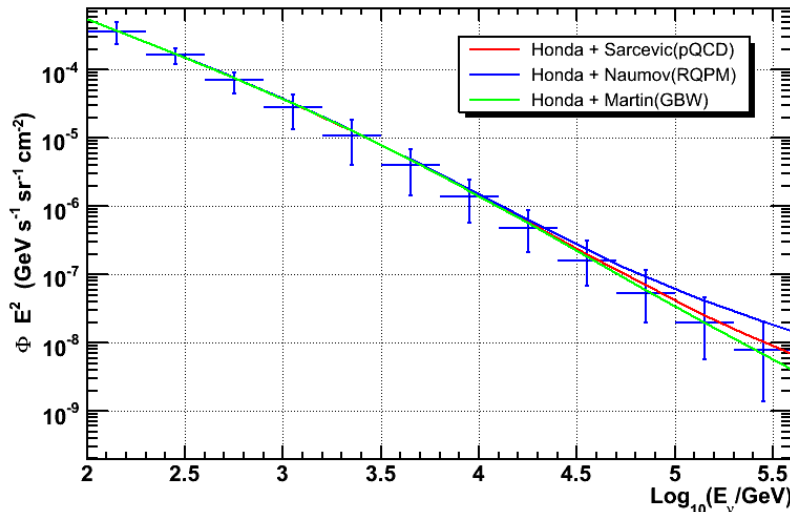
Used only 97-180 zenith region in unfolding: 17689 events from IC40 data set
Observable was the reconstructed muon dE/dX



$$b(dE/dX) = A(E_\nu, dE/dX) \Phi(E_\nu)$$



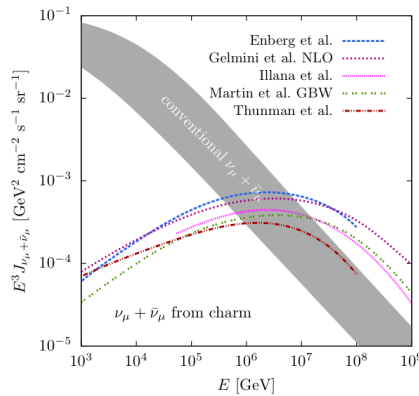
$$\Phi(E_\nu) = A^{-1}(E_\nu, dE/dX) b(dE/dX)$$



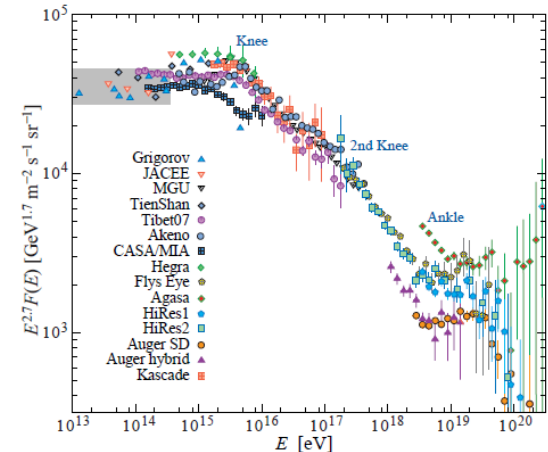
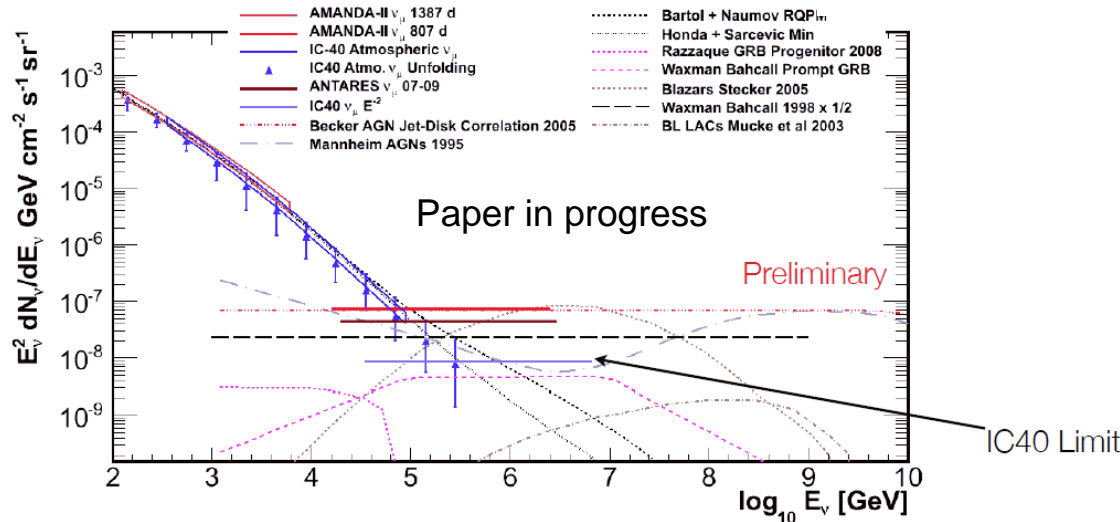
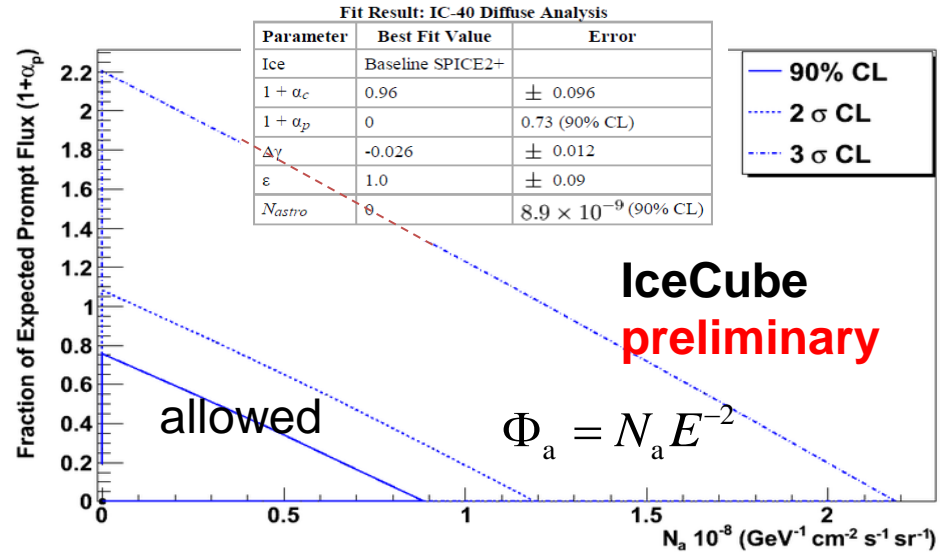
Phys. Rev. D **83**, 012001 (2011)

Prompt atmo and diffuse astrophysical

$$\Phi_p = (1 + \alpha_p) \left(\frac{E}{E_{\text{median},p}} \right)^{\Delta\gamma} \Phi_{\text{Enberg}}$$



Result consistent with only K, π atmospheric ν to 100 TeV
Charm component not yet seen; “intrinsic” charm ruled out?

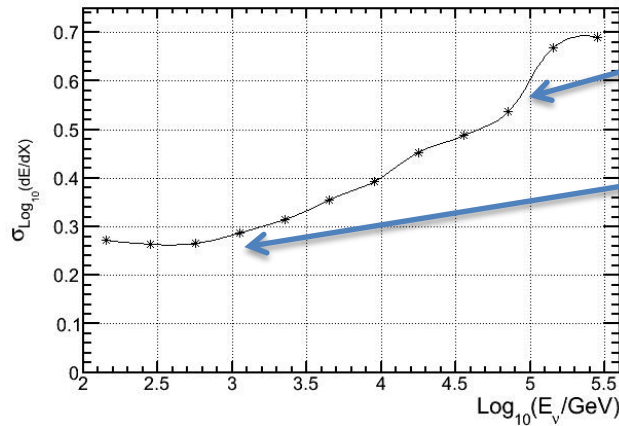
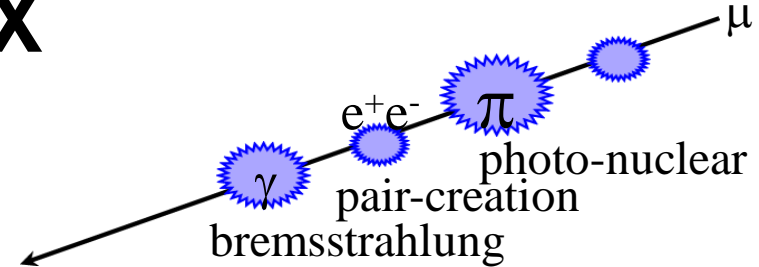


limit on astrophysical $\nu_\mu + \bar{\nu}_\mu$, assuming E^{-2} spectrum

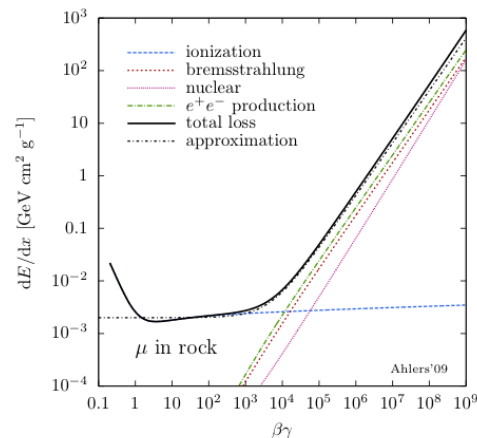
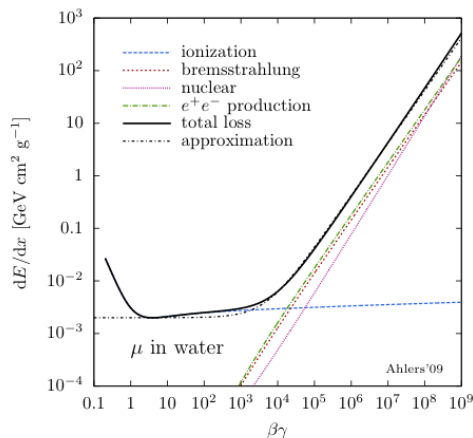
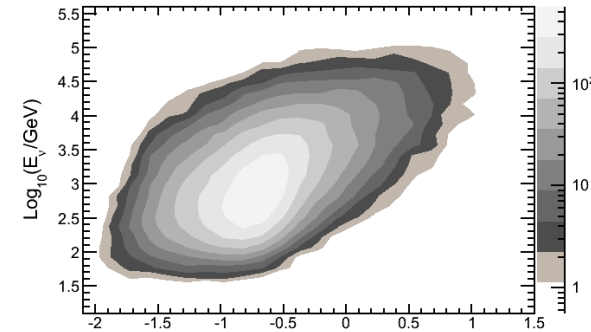
C. Amsler et al. (Particle Data Group),
Physics Letters B667, 1 (2008)

Energy Proxy: muon dE/dX

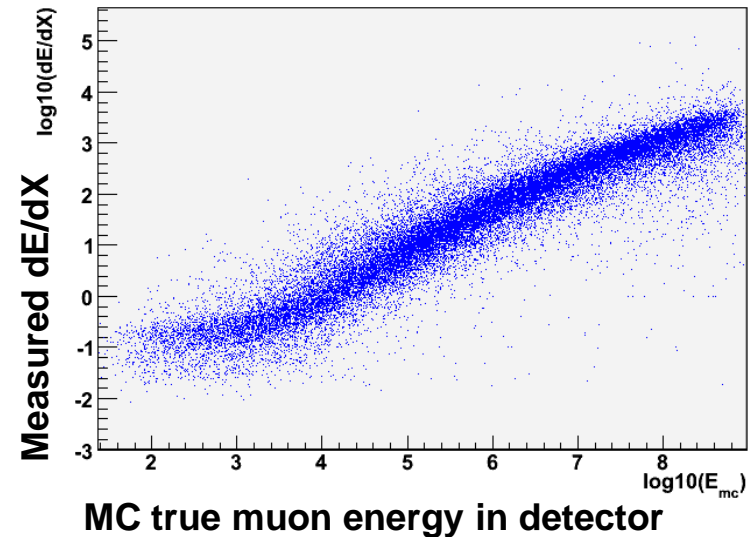
$$\underline{-\frac{dE_\mu}{dX} \approx \alpha(E_\mu) + \beta(E_\mu) E_\mu \Rightarrow \frac{\Delta E}{\text{IceCube}} \Rightarrow \frac{\sum \text{photons detected}}{\sum \text{photons predicted}}}$$



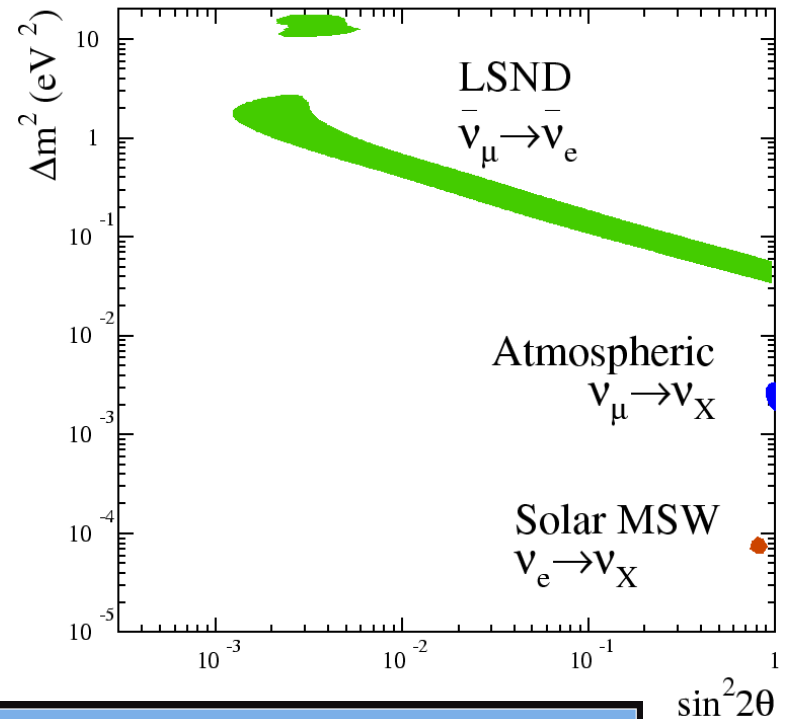
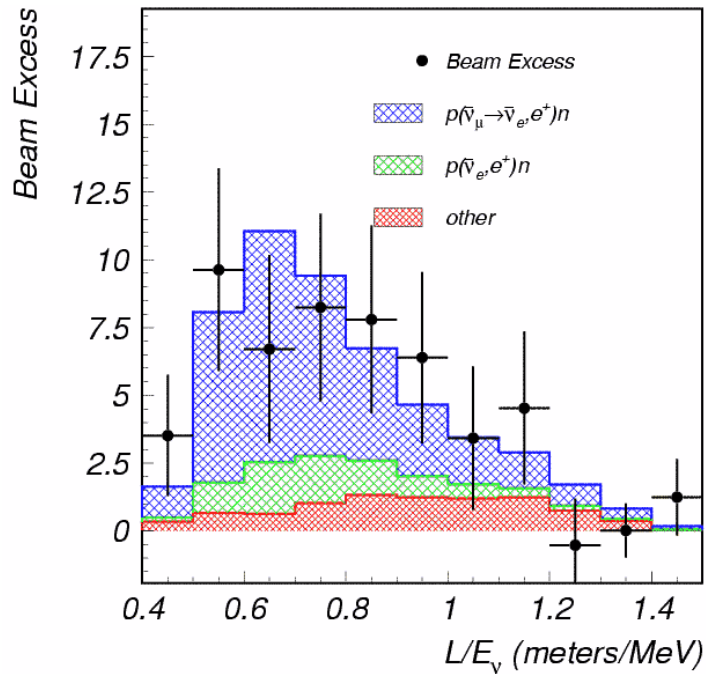
- Muon energy loss outside detector hurts correlation with neutrino energy
- Energy dependence of muon track length in detector aids energy resolution



Energy Correlation



Motivation for MiniBooNE: The LSND Evidence for Oscillations



LSND Saw an excess of $\bar{\nu}_e$:
 $87.9 \pm 22.4 \pm 6.0$ events.

With an oscillation probability of
 $(0.264 \pm 0.067 \pm 0.045)\%$.

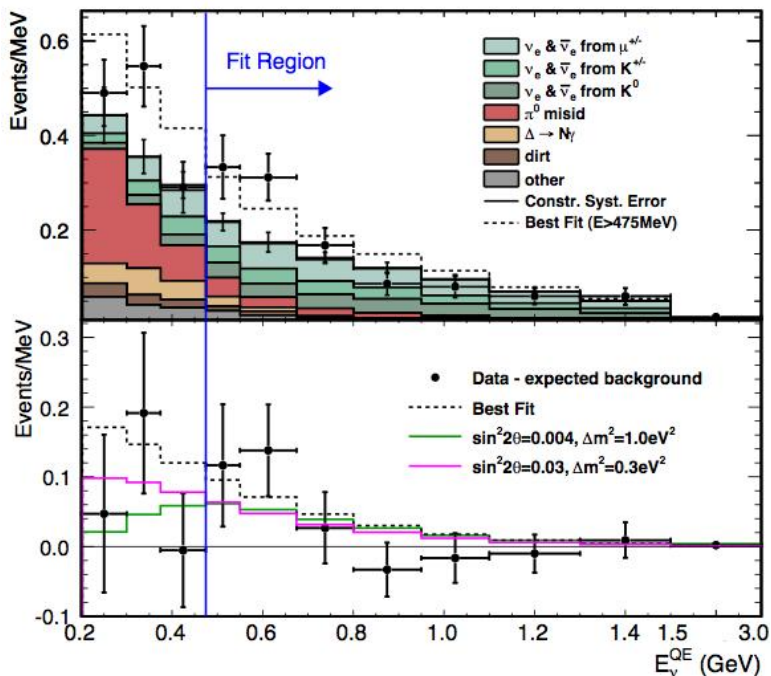
3.8 σ evidence for oscillation.

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 (eV^2) L (km)}{E (GeV)} \right)$$

The three oscillation signals cannot be reconciled without introducing Beyond Standard Model Physics!

MiniBooNE Antineutrino Results (above 475 MeV)

anti- $\bar{\nu}_e$ results

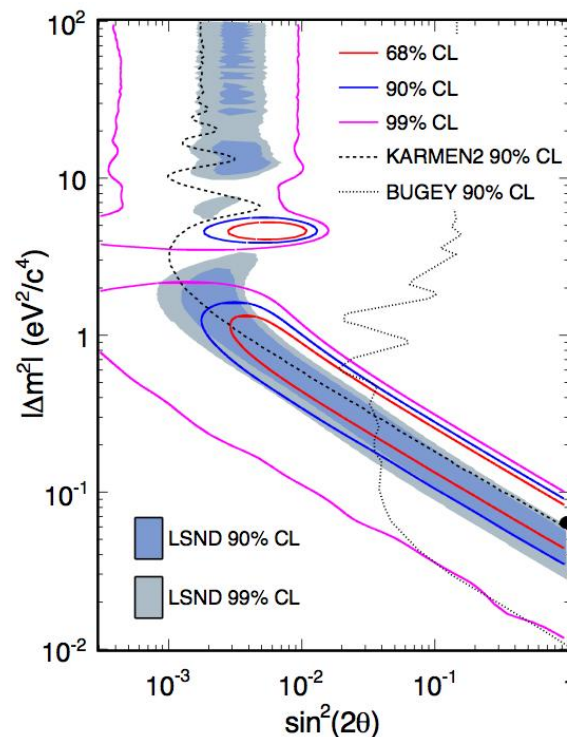
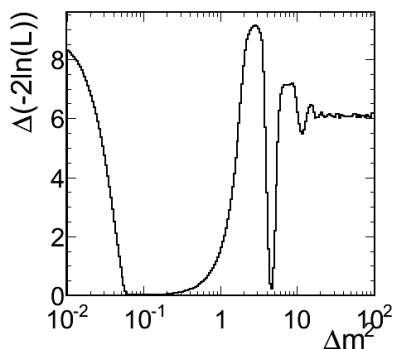


Above 475 MeV...

- ➔ In 475-1250 MeV, excess 20.9 ± 14 events (1.4σ)
- ➔ In 475-675 MeV, excess is 25.7 ± 7.2 events (3.6σ)
- ➔ True significance comes from fit over entire > 475 MeV energy region + numu constraint
- ➔ Best fit preferred over null at 99.4% CL (2.7σ)
- ➔ Probability of null hypothesis (no model dep.) is 0.5% in 475-1250 MeV signal region

Phys. Rev. Lett. **105**, 181801 (2010)

$\Delta\chi^2$ vs Δm^2

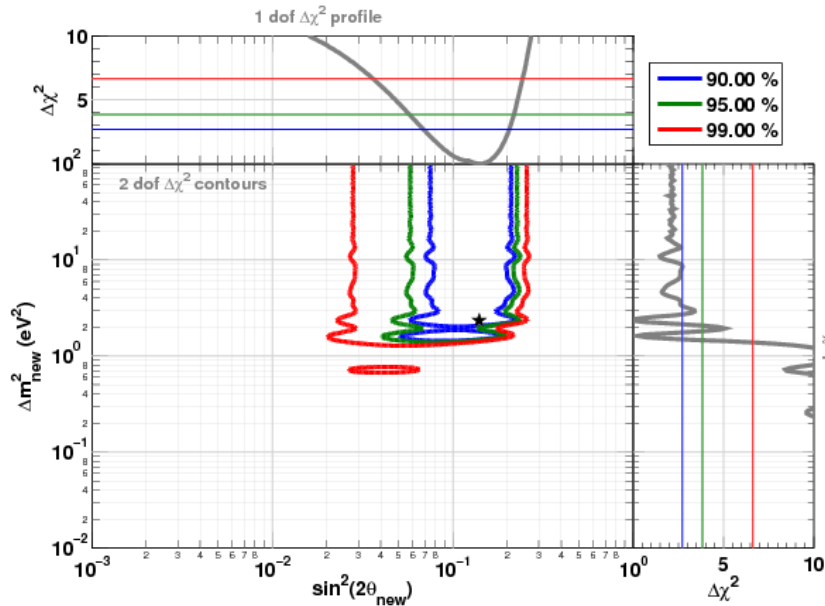


Sterile Neutrinos

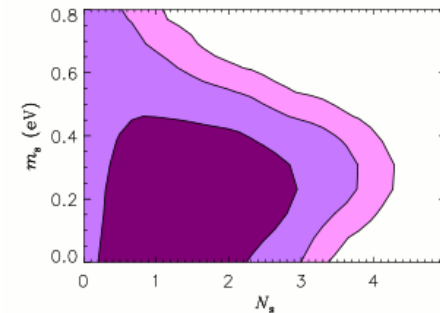
Reactor Neutrino Anomaly

(G. Mention, et. al., arXiv:1101.2755)

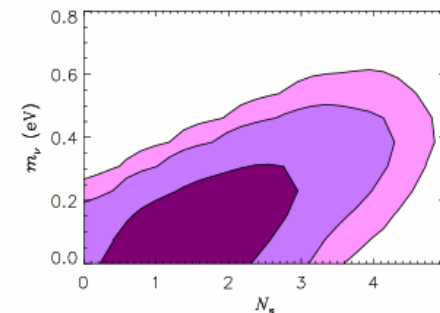
“The no-oscillation hypothesis is disfavored at 99.86% CL”



Cosmology Data Consistent with Extra, Sterile Neutrinos (J. Hamann, et. al. PRL 105 (2010) 181301)

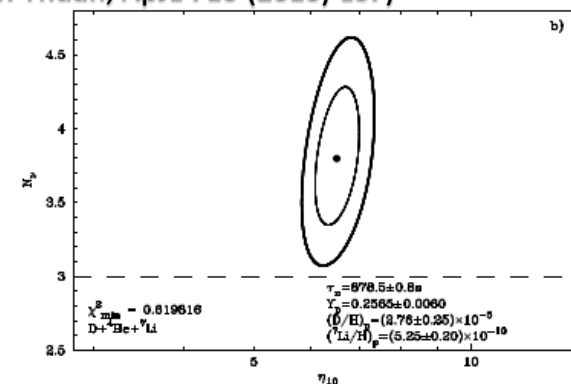
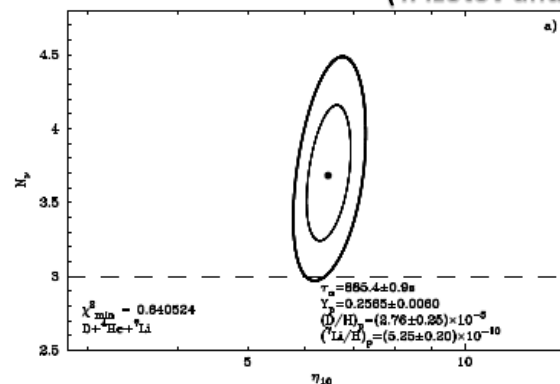


$3 + N_s$
 $m_\nu = 0$



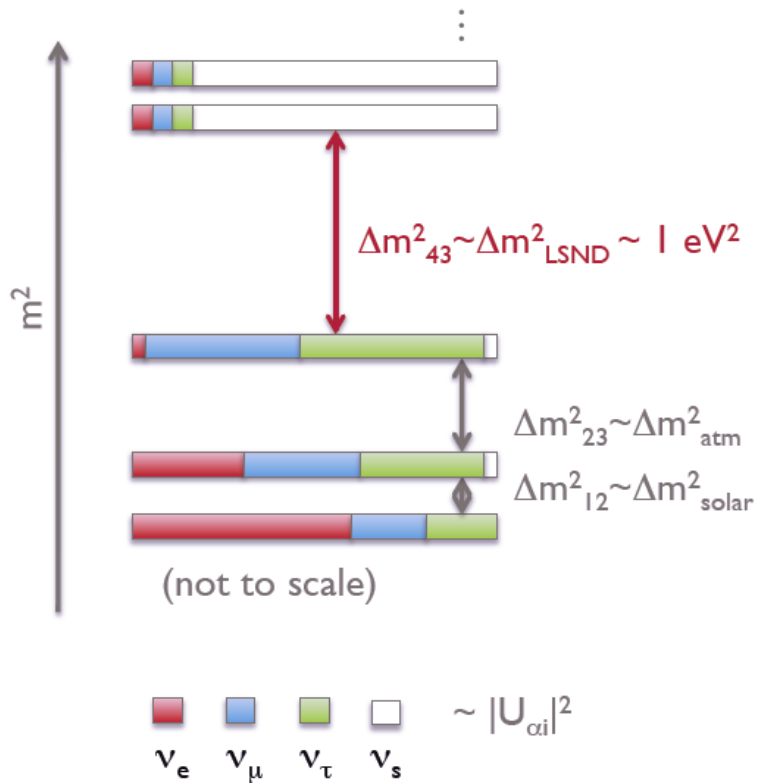
$3 + N_s$
 $m_s = 0$

(Y. Izotov and T. Thuan, ApJL 710 (2010) L67)



Mixing with sterile neutrinos?

3 active + n sterile neutrino states



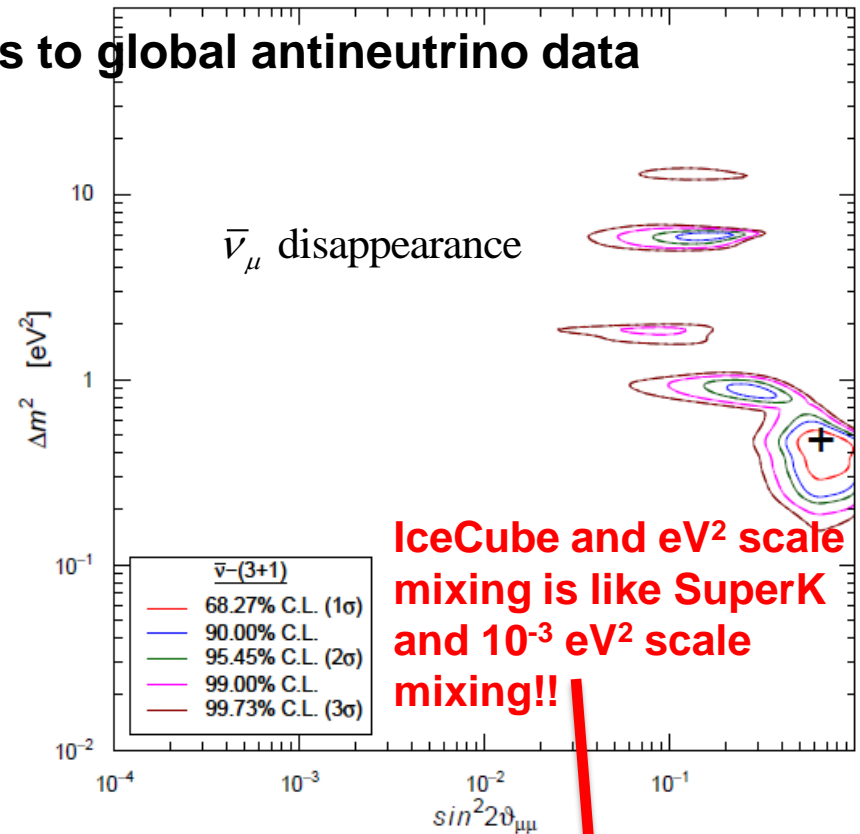
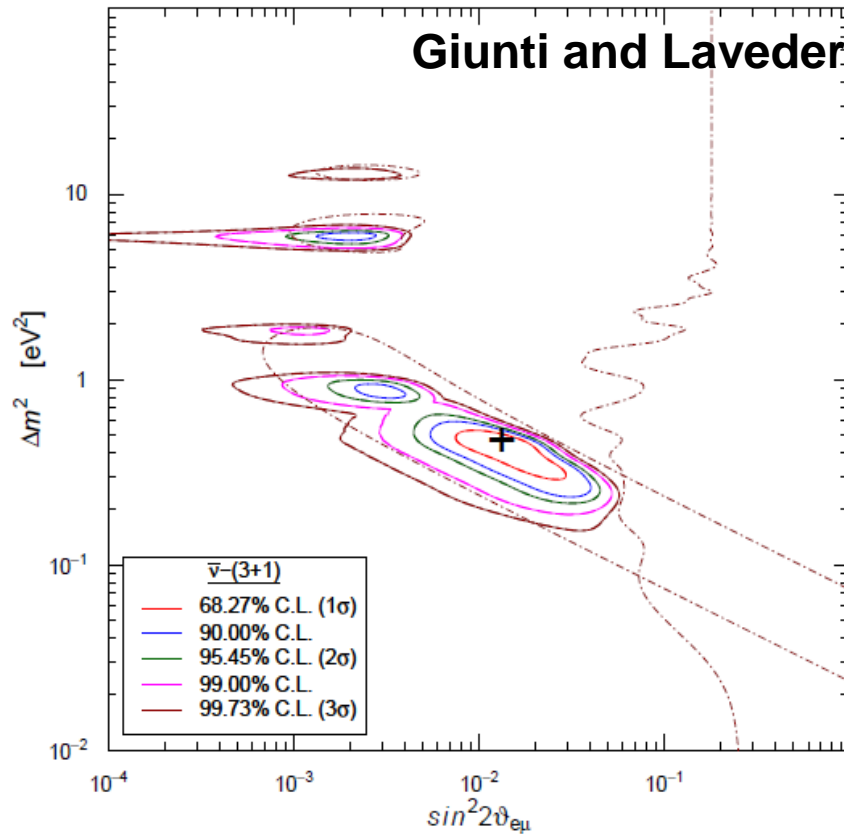
Probability of Neutrino Oscillations

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_i \sum_j |U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j}| \sin^2(1.27 \Delta m_{ij}^2 L / E_\nu)$$

As N increases, the formalism gets rapidly more complicated!

N	# Δm_{ij}^2	# θ_{ij}	#CP Phases
2	1	1	0
3	2	3	1
6	5	15	10

Antineutrino Oscillations in 3+1 Schemes



IceCube and eV² scale mixing is like SuperK and 10⁻³ eV² scale mixing!!

$$\chi^2_{\min} = 82.0 \quad \text{NdF} = 83 \quad \text{GoF} = 51\%$$

$$\Delta m^2 = 0.45 \text{ eV}^2 \quad \sin^2 2\vartheta_{e\mu} = 0.013 \quad \sin^2 2\vartheta_{ee} = 0.017 \quad \sin^2 2\vartheta_{\mu\mu} = 0.65$$

Prediction: large SBL $\bar{\nu}_\mu$ disappearance at $0.1 \lesssim \Delta m^2 \lesssim 1 \text{ eV}^2$

[Giunti, Laveder, arXiv:1012.0267]

Matter Effects for $\sim 1\text{eV}^2$ scale oscillations

S. Choubey, arXiv:0709.1937

$$\Delta m_M^2 = \sqrt{(\Delta m_{vac}^2 \cos 2\theta_{vac} - A_M)^2 + (\Delta m_{vac}^2 \sin 2\theta_{vac})^2}$$

$$\sin 2\theta_M = \frac{\Delta m_{vac}^2 \sin 2\theta_{vac}}{\Delta m_M^2}$$

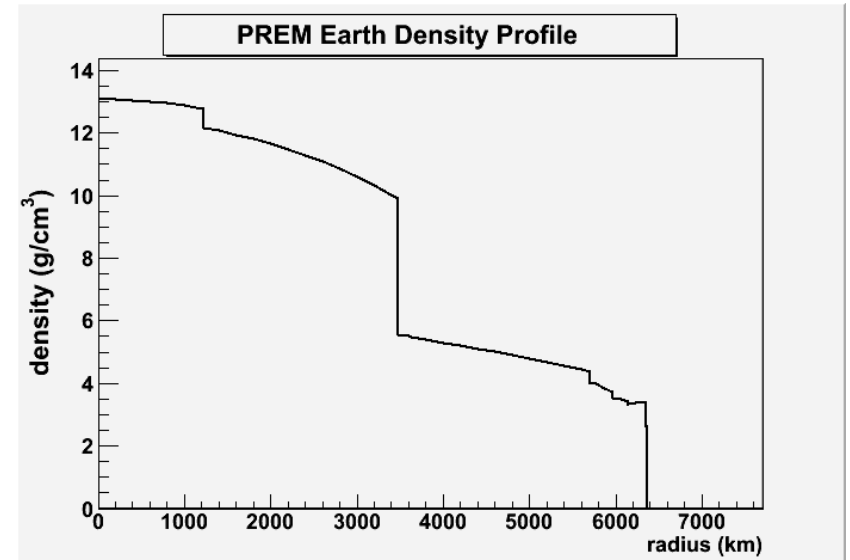
$$\cos 2\theta_M = \frac{\Delta m_{vac}^2 \cos 2\theta_{vac} - A_M}{\Delta m_M^2}$$

Due to $\sim 1\text{eV}^2$ mass scale, density of Earth, and the fact that the mixing is with sterile neutrinos, can't ignore "MSW" effect. Effective mixing parameters are now functions of energy and zenith angle

$$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_s : A_M = A_{NC} = \sqrt{2} E_\nu G_F N_n$$

$$\text{for } \nu_\mu \leftrightarrow \nu_s : A_M = -A_{NC}$$

$$N_n \approx N_e \approx 0.5 N_A \rho$$



Matter Effects for $\sim 1\text{eV}^2$ scale oscillations

$$1 - \sin^2(2\theta_M) \sin^2\left(1.27 \frac{\Delta m_M^2 (\text{eV}^2) L(\text{km})}{E(\text{GeV})}\right)$$

$$i \frac{d}{dx} \Psi = \frac{1}{2E} (UM^2U^\dagger + A) \Psi$$

$$i \frac{d}{dx} \begin{pmatrix} \psi_{\mu\mu} \\ \psi_{\mu s} \end{pmatrix} = \frac{\Delta m_M^2}{4E_\nu} \begin{pmatrix} -\cos 2\theta_M & \sin 2\theta_M \\ \sin 2\theta_M & \cos 2\theta_M \end{pmatrix} \begin{pmatrix} \psi_{\mu\mu} \\ \psi_{\mu s} \end{pmatrix}$$

$$P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} = |\psi_{\mu\mu}|^2$$

Can generalize to include more than one active and/or more than one sterile neutrino

C. Giunti and C.W. Kim, *Fundamentals of Neutrino Physics and Astrophysics*, Oxford University Press, (2007)

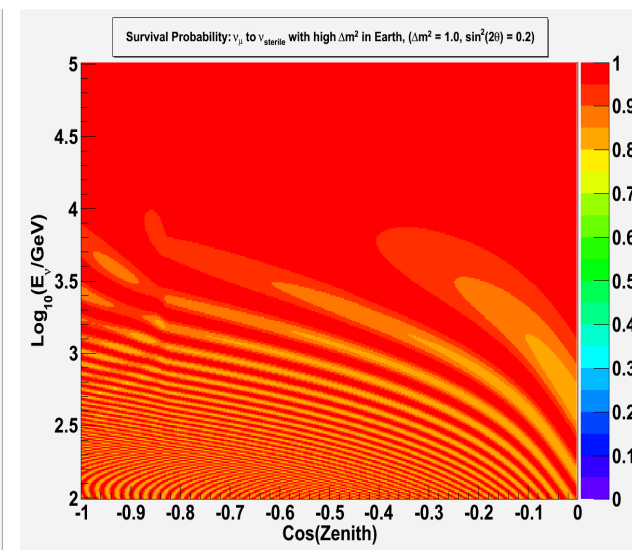
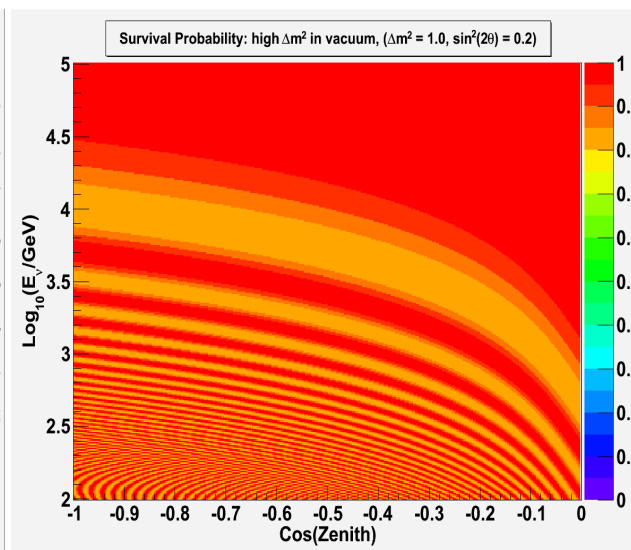
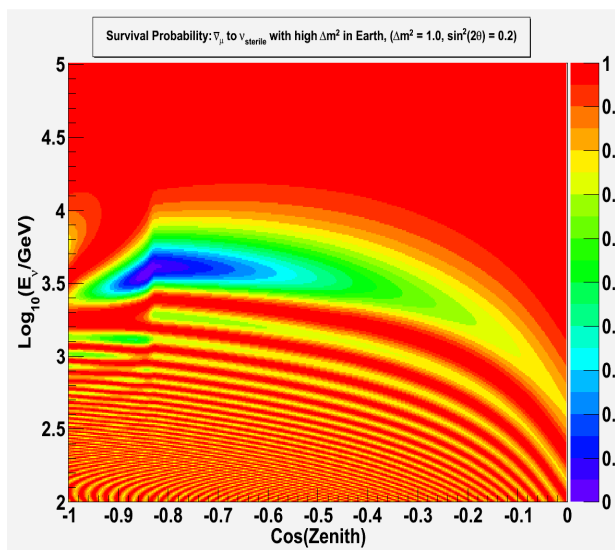
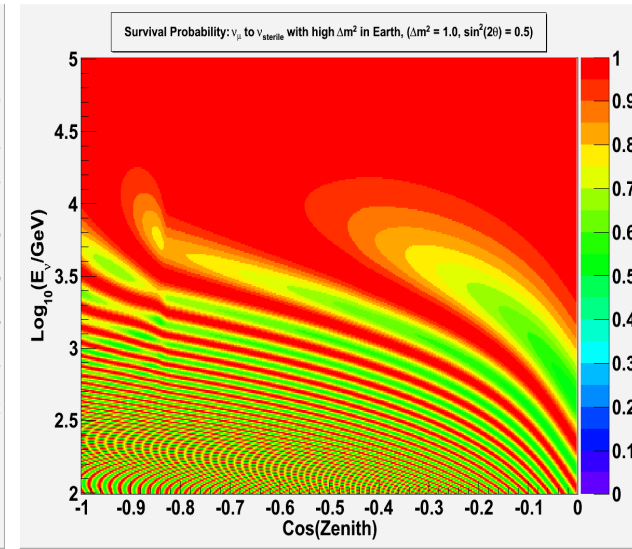
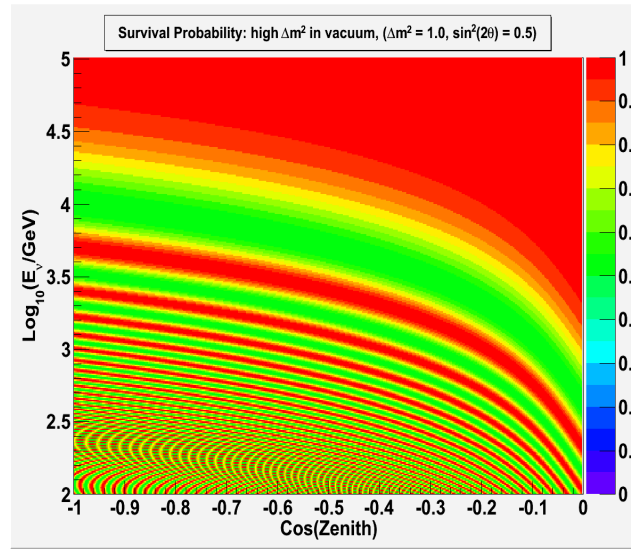
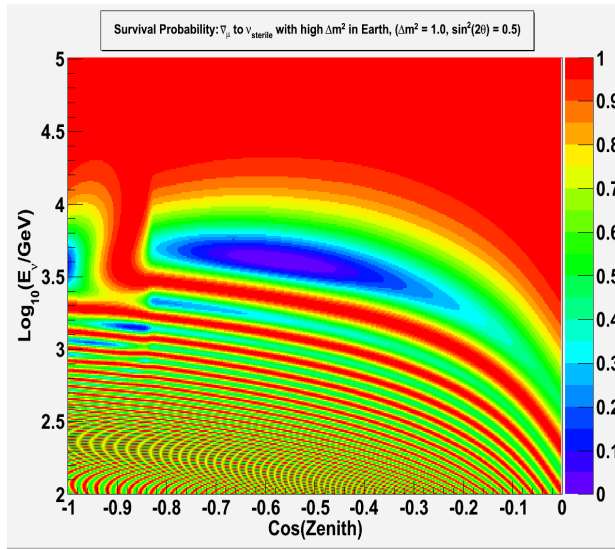
Matter Effects in IceCube

Oscillation survival probability as function of neutrino energy and zenith

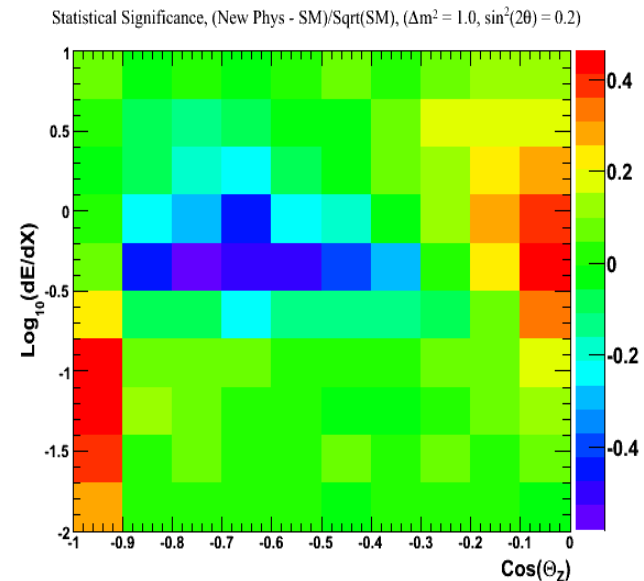
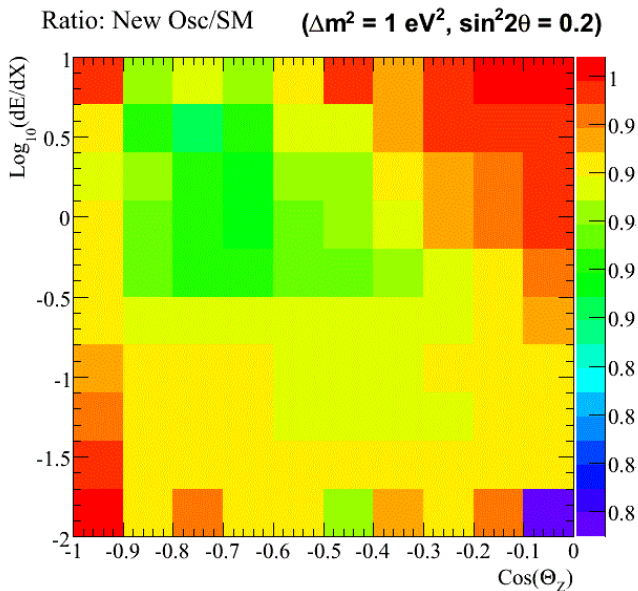
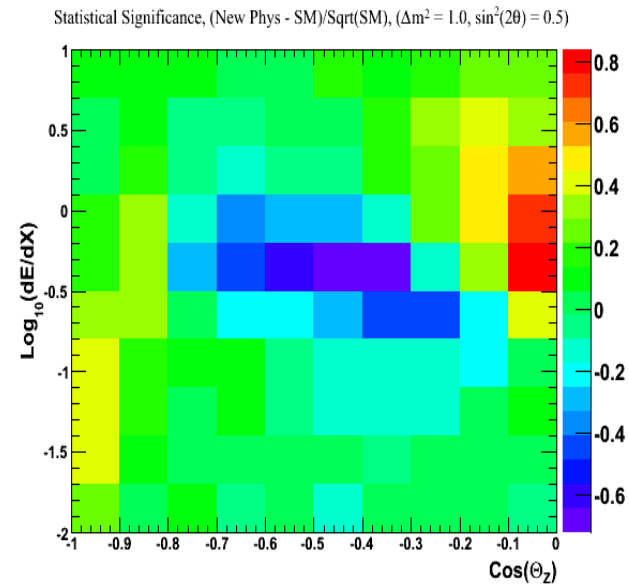
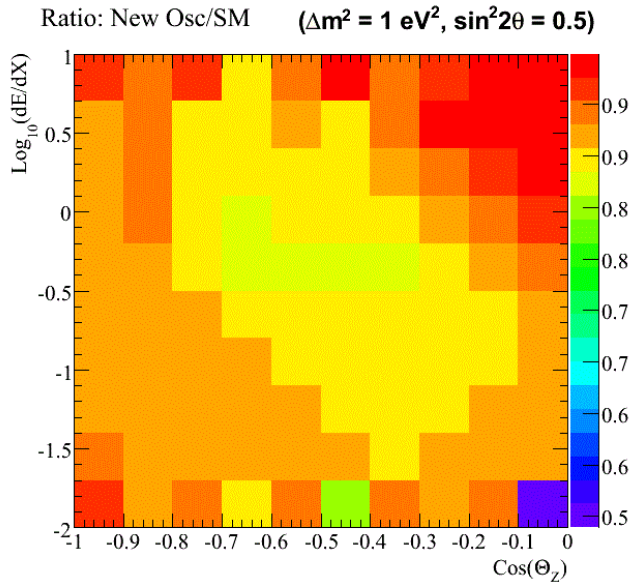
$\bar{\nu}_\mu$ to sterile in Earth

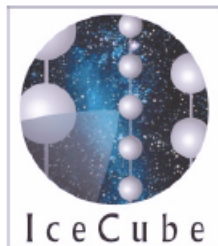
ν_μ or $\bar{\nu}_\mu$ to sterile in vacuum

ν_μ to sterile in Earth



$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_s$ Signatures in 40-str IceCube





Methodology: Likelihood Analysis

Likelihood ratio with Feldman-Cousins ordering principle to determine exclusion regions

Binned in Energy Estimator (dE/dX) and $\cos(\theta_z)$

- 10 bins in dE/dX
- 10 bins in $\cos(\theta_z)$

$$R \sim \chi^2(\theta_r^{\text{current point}}, \theta_s^{\text{for BF at current pt}}) - \chi^2(\theta_r^{\text{BF point}}, \theta_s^{\text{BF point}})$$

Test statistic:

$$R = -2 \ln \frac{L_0}{\hat{L}}$$

L_0 = max likelihood with physics params fixed (nuisance params free)

\hat{L} = max likelihood with all params free

Likelihood function:

$$L(\{n_{ij}\} | \{\mu_{ij}(\theta_r, \theta_s)\}) = \prod_{i,j} \frac{\mu_{ij}^{n_{ij}}}{n_{ij}!} e^{-\mu_{ij}}$$

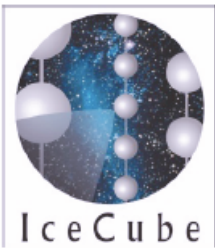
θ_r : Physics Parameters

θ_s : Nuisance Parameters

Feldman-Cousins ordering:

- (1) compute R 's for ~ 1000 toy MC experiments at each point in physics parameter space
- (2) R_{crit} = value of R such that $X\%$ of MC experiments (performed at that point in physics parameter space) have $R < R_{crit}$
- (3) point is within allowed region if $R_{data} < R_{crit}$ at that point

Nuisance parameters handled using profile construction method



Flux Modeling

Event weights in predicted flux:

$$w = \epsilon \left(\frac{E}{E_{ref}} \right)^{\Delta\gamma} \left[1 + 2\alpha \left(\cos\theta_z + \frac{1}{2} \right) \right] \left\{ \left[1 + 2\alpha_c \left(\cos\theta_z + \frac{1}{2} \right) \right] w_C + A_p \left(\frac{E}{E_{ref,p}} \right)^{\Delta\gamma_p} w_P \right\} P_{\nu_\mu \rightarrow \nu_\mu}(\theta_r)$$

w_C : "Honda_2006" conventional atmospheric neutrino spectrum
(Honda et al, PRD 75, 043006 (2007))

$$E_{ref} \approx 1 \text{ TeV}$$

w_P : "Sarcevic_std" prompt atmospheric neutrino spectrum
(Enberg, et al, hep-ph 0806.0418)

$$E_{ref,p} \approx 5 \text{ TeV}$$

(atmospheric neutrino spectra extended to energies above a TeV in NeutrinoFlux Class)

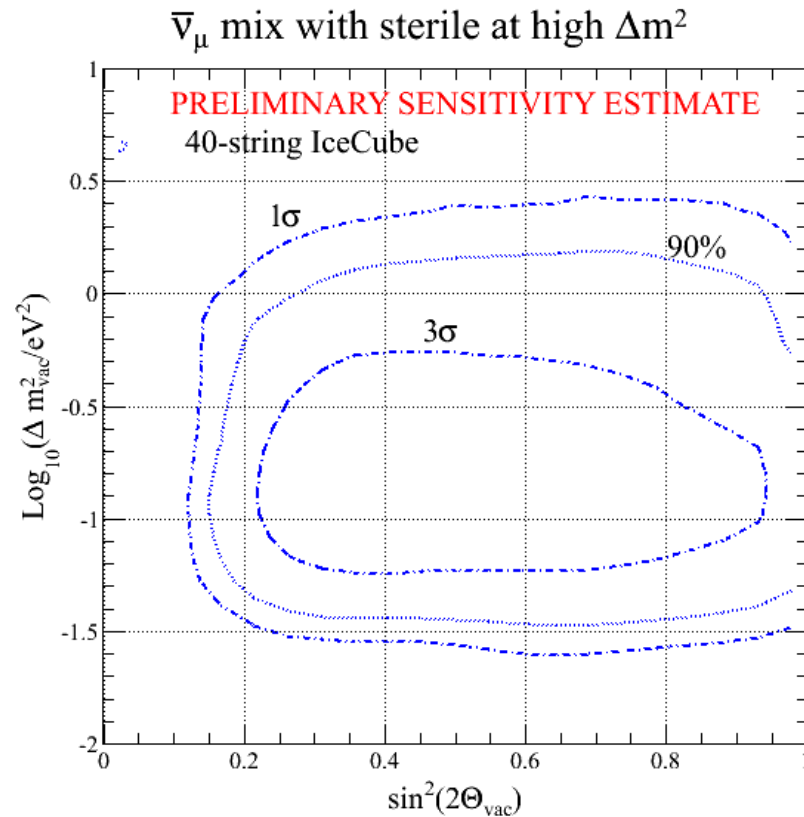
$P_{\nu_\mu \rightarrow \nu_\mu}(\theta_r)$: oscillation or decoherence survival probability;
 θ_r = model dependent physics parameters

Nuisance Parameters (θ_s): $\epsilon, \Delta\gamma, \alpha, \alpha_c, A_p, \Delta\gamma_p$

Muons from taus:

- If muon neutrinos are oscillating to tau neutrinos, some of these tau neutrinos can produce muons in detector via tau's decay chain
- Can account for in simulation chain

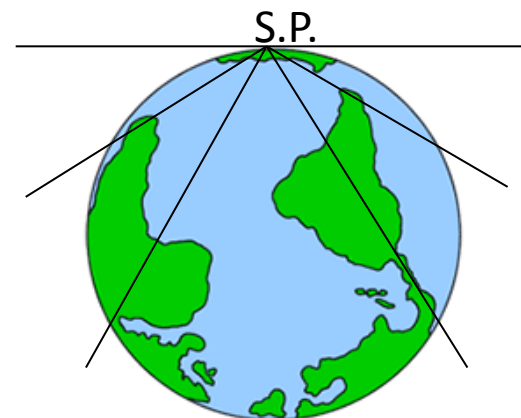
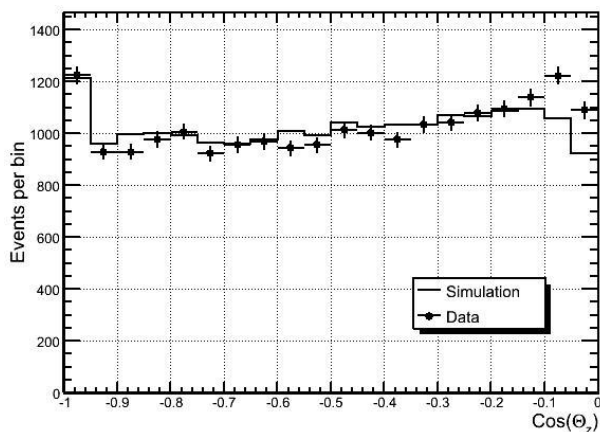
IC40 sensitivity for sterile neutrinos



plot assumes only $\bar{\nu}_\mu \leftrightarrow \bar{\nu}_s$

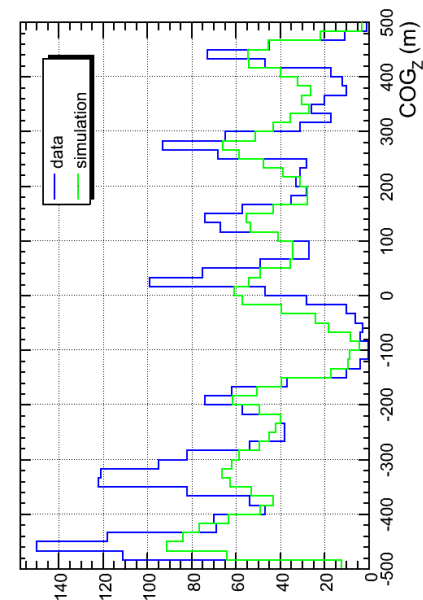
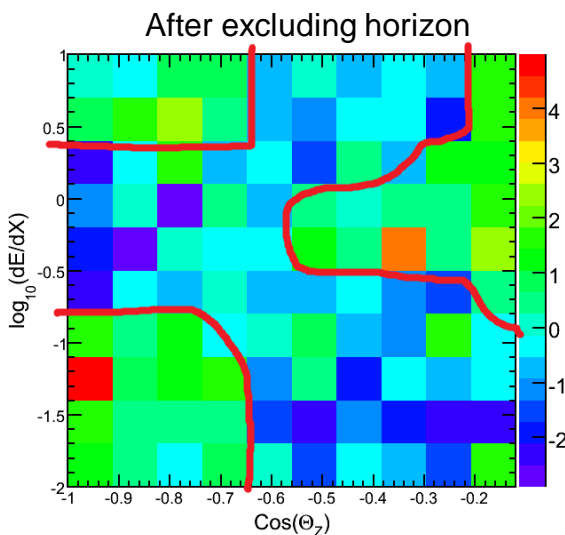
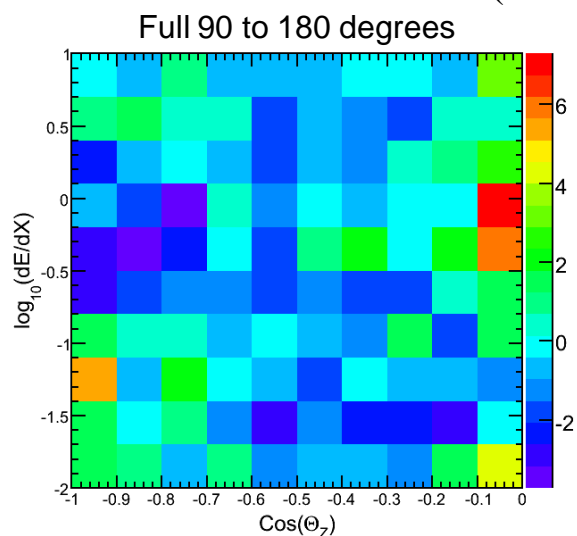
Based on likelihood analysis with nuisance parameters

Apparent “Excess” near horizon



Bin-by-bin statistical significance of data/MC disagreement

$$(data - mc) / \sqrt{mc}$$



Systematic Uncertainties

There are known knowns; there are things we know we know.

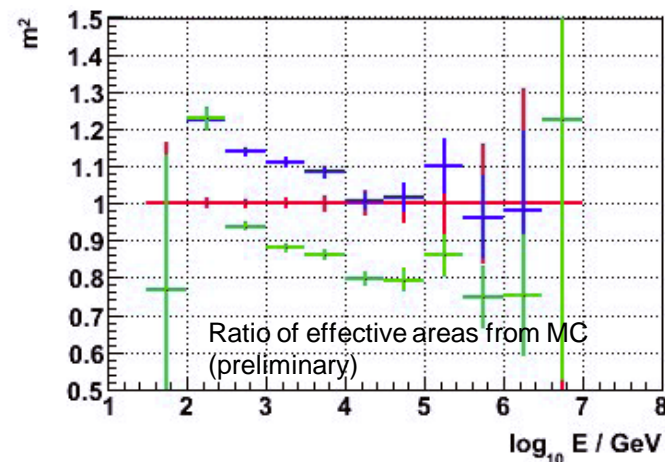
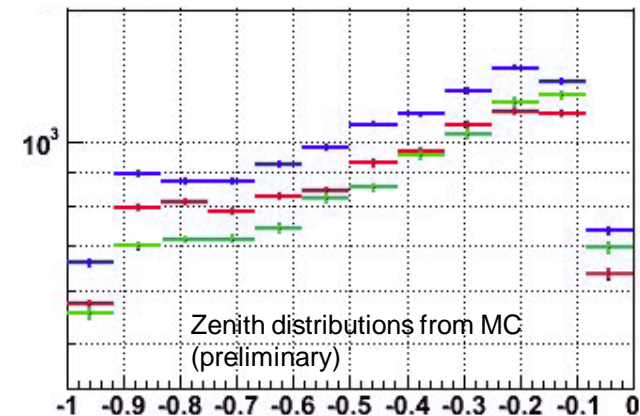
We also know there are known unknowns; that is to say we know there are some things we do not know.

But there are also unknown unknowns, the ones we don't know we don't know.

– Donald Rumsfeld

- Atmospheric variability?
- Ice Model/photon prop code?
- Cosmic ray muon background?
- Muon propagation code?
- DOM sensitivity and angular acceptance?
- Hole ice?
- New physics?
- Other?

Impact of systematic uncertainties in ice modeling:



Alternate test statistic with error matrix of correlated errors

(Nuisance parameters do not adequately describe energy and zenith dependent uncertainties and correlations between bins)

$$\chi^2 = \sum_{i=1}^N \sum_{j=1}^N (X_i - D_i) (M^{-1})_{ij} (X_j - D_j) + \ln(|M|)$$

N = number of bins (80 for $8 \times 10 \text{ } dE / dX$ vs. $\cos(\text{zen})$)

M has 6400 elements

$$M_{ij,\alpha} = (X_i(\alpha) - X_i(CV))(X_j(\alpha) - X_j(CV))$$

α is a particular systematic variation, CV is central value

$$M_{ij} = \sum_{\alpha} M_{ij,\alpha} + \delta_{ij} x_i$$

$M_{ij,\alpha}$: OM sensitivity, ice model, photon propagation, cross sections, atmospheric variability, neutrino production model, muon propagation code, cosmic ray flux (composition and spectrum), reconstructions, angular acceptance, etc...

Alternate test statistic with error matrix of correlated errors

$$\chi^2 = \sum_{i=1}^N \sum_{j=1}^N (X_i - D_i) (M^{-1})_{ij} (X_j - D_j) + \ln(|M|)$$

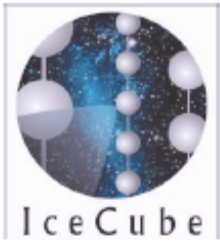
Test Statistic:

$$\Delta\chi^2 = \chi^2(X(\theta), M(\theta)) - \chi^2(X(\theta_{bf}), M(\theta_{bf}))$$

where θ represents the physics parameters for the disappearance model being tested and bf is the best fit point

Can also include nuisance parameters with pull terms:

$$\frac{(\bar{\alpha} - \alpha(\theta))^2}{\sigma_{\alpha}^2}, \quad \bar{\alpha} = \text{expected mean for param } \alpha$$



Summary

- IceCube is fully functional and searching for a variety of anticipated and unknown astrophysical signals
 - Anticipate new discoveries in next couple(few) years
- The characteristics of IceCube that will enable it to open up a neutrino window to the universe also make it ideal for studying high energy atmospheric neutrinos
- IceCube will be able to weigh in on the sterile neutrino debate, as well as other new physics models that alter the energy and zenith dependence of the flux of high energy atmospheric neutrinos
- Once we get the systematic uncertainties under control...
- Plan to perform the analysis on IC59 data